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Electric potential
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ELECTRIC POTENTIAL(V)

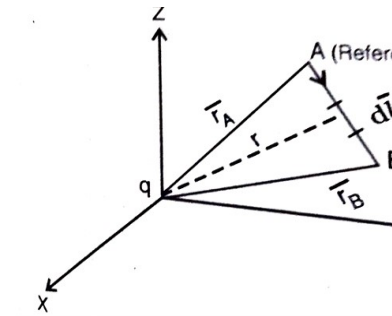
Electric potential (V) is defined as the work done in moving a unit positive charge from a reference point to a point in an electric field.

ELECTRIC FIELD POTENTIAL DUE TO POINT CHARGE

In the case of electric field the reference point is taken as infinity.

Field produced by a charge q at a distance r is,

$$\vec{E} = \frac{q}{4\pi\epsilon_0 r^2} \hat{r}$$



When a unit positive charge is moved from A to B as shown in the figure, work is to be done against the field.

The small work to be done to move unit positive charge for a very small displacement dl ,

$$dw = F \cdot dl$$

$F = -E$, which is the force applied to counter the field.

Total work to be done in moving unit positive charge from a reference point to a given point in the field,

$$\begin{aligned}
 & \int_A^B dw \\
 & = -\int_A^B \bar{E} \cdot d\bar{l} \\
 & = -\int_A^B \bar{E} \cdot d\bar{l}
 \end{aligned}$$

shows that work done can be obtained as the negative line integral of the electric field. As per the definition of electric potential, V ,

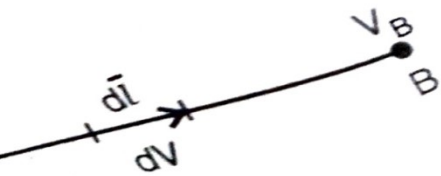
$$\begin{aligned}
 V &= w = - \int_{A=\infty}^{B=(\text{point m } \bar{E})} \bar{E} \cdot d\bar{l} \\
 d\bar{l} &= dr\hat{r} + rd\theta\hat{\theta} + r \sin\theta d\phi\hat{\phi} \\
 \bar{E} \cdot d\bar{l} &= \frac{q}{4\pi\epsilon_0 r^2} \hat{r} \cdot (dr\hat{r} + rd\theta\hat{\theta} + r \sin\theta d\phi\hat{\phi}) \\
 &= \frac{q}{4\pi\epsilon_0 r^2} dr \\
 \therefore V &= - \int_{\infty}^{r=r_B} \frac{q}{4\pi\epsilon_0 r^2} dr \\
 V &= \frac{-q}{4\pi\epsilon_0} \int_{\infty}^{r_B} \frac{dr}{r^2} \\
 &= \frac{-q}{4\pi\epsilon_0} \left[-\frac{1}{r} \right]_{\infty}^{r_B} = \frac{q}{4\pi\epsilon_0} \left[\frac{1}{r_B} - \frac{1}{\infty} \right] \\
 V &= \frac{q}{4\pi\epsilon_0 r_B} \quad \text{-----(7)}
 \end{aligned}$$

Eqn. (7) gives the expression for potential at the point B with position vector r_B . In general,

$$V = \frac{q}{4\pi\epsilon_0 r}$$

POTENTIAL DIFFERENCE

Potential difference is defined as the work done on a unit positive charge between two points. If a unit positive charge is moved from A to B, then potential difference = $V_B - V_A$, where V_B and V_A are the electric potential at the point B and A respectively. Here $V_B > V_A$ since work is done on the unit positive charge to move from A to B. If dV represents a very small potential difference along the path connecting A and then,



$$V_B - V_A = \int_A^B dV$$
$$dV = -\vec{E} \cdot d\vec{l}$$

$$\begin{aligned}\therefore V_B - V_A &= - \int_A^B \vec{E} \cdot d\vec{l} \\ &= - \int_{r_A}^{r_B} \frac{q}{4\pi\epsilon_0 r^2} \hat{r} \cdot (dr\hat{r} + r d\theta\hat{\theta} + r \sin\theta d\phi\hat{\phi}) \\ &= - \int_{r_A}^{r_B} \frac{q}{4\pi\epsilon_0 r^2} dr \\ &= - \frac{q}{4\pi\epsilon_0} \int_{r_A}^{r_B} \frac{dr}{r^2} = - \frac{q}{4\pi\epsilon_0} \left[-\frac{1}{r} \right]_{r_A}^{r_B} \\ V_B - V_A &= \frac{q}{4\pi\epsilon_0} \left[\frac{1}{r_B} - \frac{1}{r_A} \right]\end{aligned}$$

ELECTRIC FIELD AS THE NEGATIVE GRADIENT OF POTENTIAL

represents a small p.d. between two close points in an electric field of intensity \vec{E} . If $d\vec{l}$ represents the distance between two points, then,

$$dV = dw = -\vec{E} \cdot d\vec{l}$$

Mathematically dV can be considered as the total differential of V .

In SPC system.
$$dV = \frac{\partial V}{\partial r} dr + \frac{\partial V}{\partial \theta} d\theta + \frac{\partial V}{\partial \phi} d\phi$$

It can be proved that r.h.s. of eqn. (2)

$$= \nabla V \cdot d\vec{l}$$

where $d\vec{l}$ is the elemental displacement in SPC.

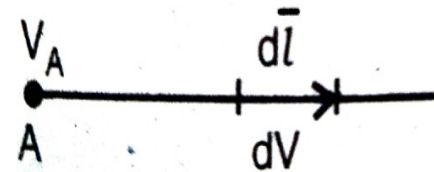
$$\text{I.e., } dV = \nabla V \cdot d\vec{l}$$

Comparing eqn. (1) and eqn. (2)

$$-\vec{E} \cdot d\vec{l} = \nabla V \cdot d\vec{l}$$

$$-\vec{E} = \nabla V$$

$$\therefore \vec{E} = -\nabla V$$



POTENTIAL DUE TO INFINITELY LONG WIRE

V be the distance r from an infinitely long line charge distribution having linear charge density λCm^{-1}
potential is given by,

$$V = - \int_{\text{ref point}}^{\text{point}} \ln E \cdot dl$$

$$\vec{E} = \frac{\lambda}{2\pi\epsilon_0 r} \hat{r}$$

$$d\vec{l} = dr\vec{r} + r d\theta\hat{\theta} + r \sin\theta d\phi\hat{\phi}$$

and reference point is given by", $r = a$

$$\therefore V = \int_{r=a}^{r=r} \frac{\lambda}{2\pi\epsilon_0 r} \hat{r} \cdot (dr\vec{r} + r d\theta\hat{\theta} + r \sin\theta d\phi\hat{\phi})$$

$$= - \int_a^r \frac{\lambda}{2\pi\epsilon_0 r} dr$$

$$= \frac{-\lambda}{2\pi\epsilon_0} \int_a^r \frac{dr}{r}$$

$$= \frac{-\lambda}{2\pi\epsilon_0} [\log r]_a^r = \frac{-\lambda}{2\pi\epsilon_0} [\log r - \log a]$$

$$V = \frac{1}{2\pi\epsilon_0} [\log a - \log r]$$

POTENTIAL V OUTSIDE THE SHELL($r>R$)

$$V_{\text{out}} = - \int_{r=\infty}^{r=r} \bar{\mathbf{E}}_{\text{out}} \cdot d\bar{\mathbf{l}}$$

$$V_{\text{out}} = - \int_{\infty}^r \frac{q}{4\pi\epsilon_0 r^2} \hat{r} \cdot (dr\hat{r} + rd\theta\hat{\theta} + r\sin\theta d\phi\hat{\phi})$$

$$V_{\text{out}} = - \int_{\infty}^r \frac{q}{4\pi\epsilon_0 r^2} dr$$

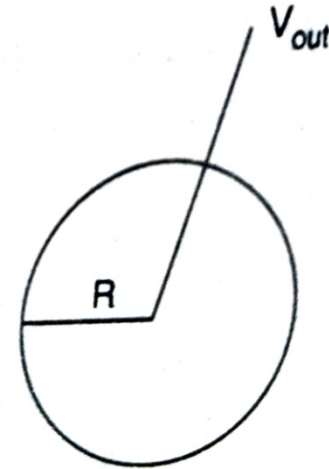
$$V_{\text{out}} = - \frac{q}{4\pi\epsilon_0} \int_{-\infty}^r \frac{1}{r^2} dr$$

$$= - \frac{q}{4\pi\epsilon_0} \left[\frac{-1}{r} \right]_{\infty}^r$$

$$= \frac{q}{4\pi\epsilon_0} \left[\frac{1}{r} - \frac{1}{\infty} \right]$$

$$V_{\text{out}} = \frac{q}{4\pi\epsilon_0 r}$$

$$q = 4\pi R^2 \sigma$$



here R , radius of the shell and σ surface charge density.

$$V_{\text{out}} = \frac{4\pi R^2 \sigma}{4\pi\epsilon_0 r}; V_{\text{out}} = \frac{R^2 \sigma}{\epsilon_0 r}$$

II) POTENTIAL ON THE SURFACE OF THE SHELL (r=R)

$$V_{\text{surface}} = - \int_{r=\infty}^{r=R} \vec{E}_{\text{out}} \cdot d\vec{l}$$
$$= - \int_{\infty}^R \frac{q}{4\pi\epsilon_0 r^2} \hat{r} \cdot (dr\hat{r} + r d\theta\hat{\theta} + r \sin\theta d\phi\hat{\phi})$$

$$= - \int_{\infty}^R \frac{q}{4\pi\epsilon_0 r^2} dr$$

$$V_{\text{surface}} = \frac{-q}{4\pi\epsilon_0} \int_{\infty}^R \frac{1}{r^2} dr$$

$$= -\frac{q}{4\pi\epsilon_0} \left[\frac{-1}{r} \right]_{\infty}^R$$

$$= \frac{q}{4\pi\epsilon_0} \left[\frac{1}{r} \right]_{\infty}^R = \frac{q}{4\pi\epsilon_0} \left[\frac{1}{R} - \frac{1}{\infty} \right]$$

$$V_{\text{surface}} = \frac{q}{4\pi\epsilon_0 R}$$

$$q = 4\pi R^2 \sigma$$

$$V_{\text{surface}} = \frac{R\sigma}{\epsilon_0}$$

(II) POTENTIAL INSIDE THE SHELL ($r < R$)

r (inside)

$$= - \int_{r=\infty} \bar{E} \cdot d\bar{l}$$

$$V_{\text{in}} = - \left[\int_{r=\infty}^{r=R} \bar{E}_{\text{out}} \cdot d\bar{l} + \int_{r=R}^{r=\text{(inside)}} \bar{E}_{\text{th}} \cdot d\bar{l} \right]$$

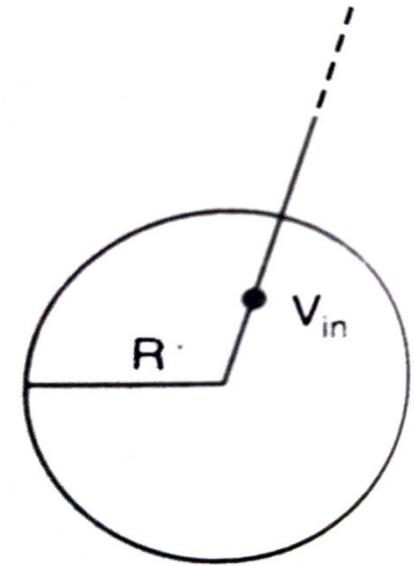
$\bar{E}_{\text{in}} =$ electric field inside the shell $= 0$

$$V_{\text{in}} = - \int_{\infty}^R \bar{E}_{\text{out}} \cdot d\bar{l}$$

$$V_{\text{in}} = \frac{q}{4\pi\epsilon_0 R} = \frac{R\sigma}{\epsilon_0}$$

$$V_{\text{in}} = V_{\text{surface}}$$

Inside the charged shell potential is constant and equal to the surface of the shell



POISSON'S EQUATION

Differential form of Gauss's law is given by

$$\nabla \cdot \bar{E} = \frac{\rho}{\epsilon_0}$$

$$\bar{E} = -\nabla V$$

$$\therefore \nabla \cdot (-\nabla V) = \frac{\rho}{\epsilon_0}$$

$$-\nabla^2 V = \frac{\rho}{\epsilon_0}$$

$$\nabla^2 V = \frac{-\rho}{\epsilon_0}$$

is called Poisson's equation.

When $\rho = 0$

$$\nabla^2 V = 0$$

is called Laplace's equation.

THANK YOU